



Direct and large-eddy simulation of rotating turbulence

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**Multiscale Modeling and Simulation (Twente)
Anisotropic Turbulence (Eindhoven)
Mathematics Department, Imperial College (London)**

**IMS Turbulence Workshop
London, March 26, 2007**

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- connect DNS to experiment - rotating decaying turbulence
- show accuracy of regularization modeling - LES
- apply regularization modeling to high Re flow
- quantify cross-over from 3D to 2D dynamics as $Ro \downarrow$
- illustrate low Ro compressibility effects

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Outline

- 1 Rotation modulates turbulence
- 2 Rotation of incompressible fluid
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Taylor-Proudman: columnar structuring

Strong rotation/small friction:

“Motion of a homogeneous fluid will be the same in all planes perpendicular to the axis of rotation”

Taylor's experiment:

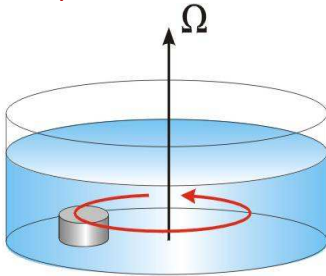
- When an object moves in a rotating flow, it drags along with it a column of fluid parallel to the rotation axis

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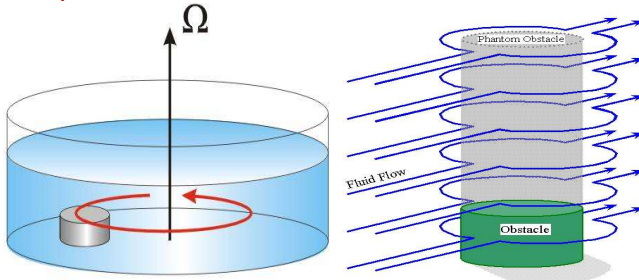
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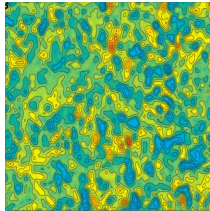
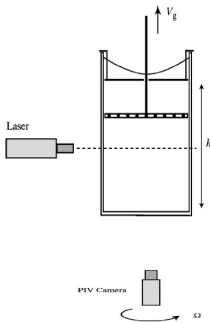
Guadalupe Islands



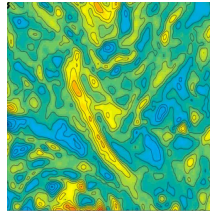
Taylor columns disturb cloud flow well above islands

Rotating turbulence in a container

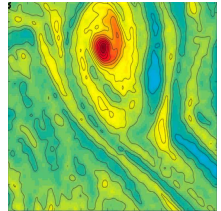
Experiment: Morize, Moisy, Rabaud, PoF, 2005



$Ro = 6.4$



$Ro = 0.6$



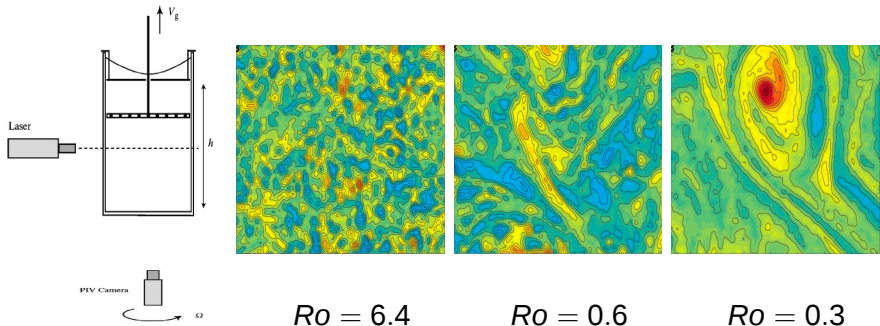
$Ro = 0.3$

Vertical structuring of flow as $Ro = u_r / (2\Omega_r \ell_r)$ decreases

Modulation: Competition 3D turbulence – 2D structuring

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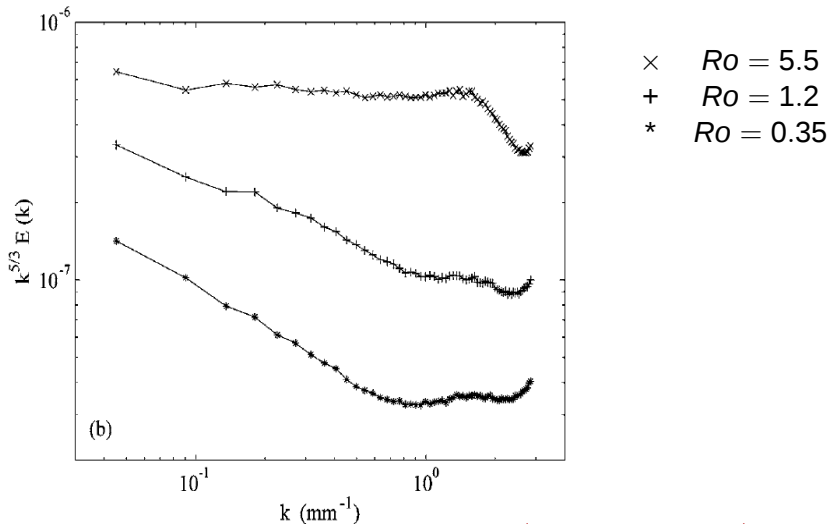
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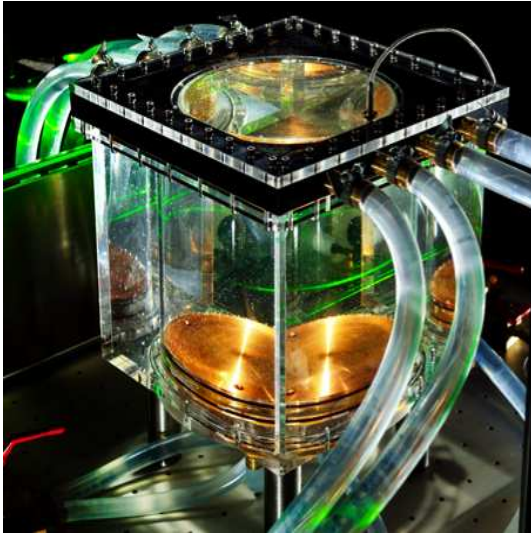
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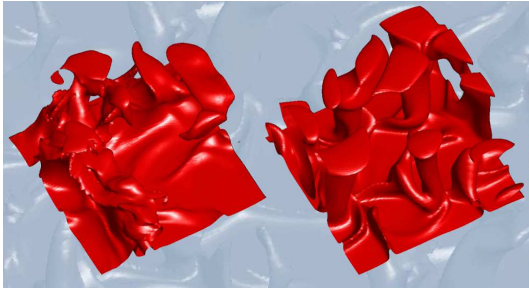
Alteration of energy-spectrum: from $k^{-5/3}$ to steeper $k^{-9/4}$

Rotating Rayleigh-Bénard problem



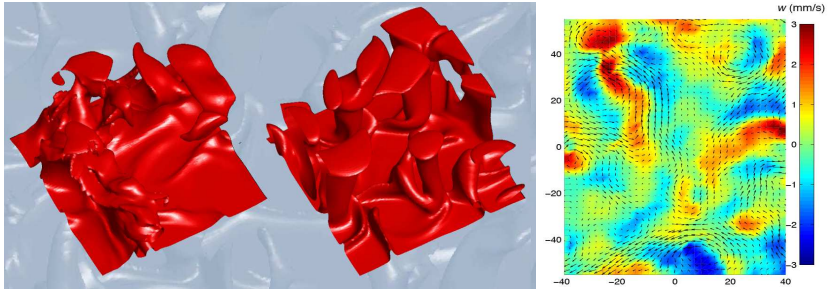
Collaboration with Rudie Kunnen, Herman Clercx - Eindhoven

Rotating Rayleigh-Bénard problem



- **Temperature:** buoyancy dominated ($Ta = 0$ - left) and rotation dominated ($Ta > Ra$)

Rotating Rayleigh-Bénard problem



- **Temperature:** buoyancy dominated ($Ta = 0$ - left) and rotation dominated ($Ta > Ra$)
- **Stereo PIV measurement of (u, v, w) in horizontal plane**

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Governing equations

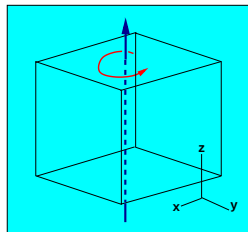
Incompressible formulation: rotating frame of reference

$$\nabla \cdot \mathbf{u} = 0$$

$$\partial_t \mathbf{u} + \nabla \cdot (\mathbf{u} \mathbf{u}) + \nabla p - \frac{1}{Re} \nabla^2 \mathbf{u} = \frac{\mathbf{u} \times \boldsymbol{\Omega}}{Ro}$$

Characteristic numbers:

$$Re = \frac{u_r \ell_r}{\nu_r} \quad ; \quad Ro = \frac{u_r}{2\Omega_r \ell_r}$$



Numerical method:

- Pseudo-spectral discretization: full de-aliasing, parallelized
- Helical-wave decomposition: Coriolis diagonalization

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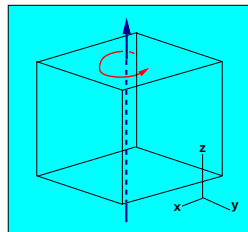
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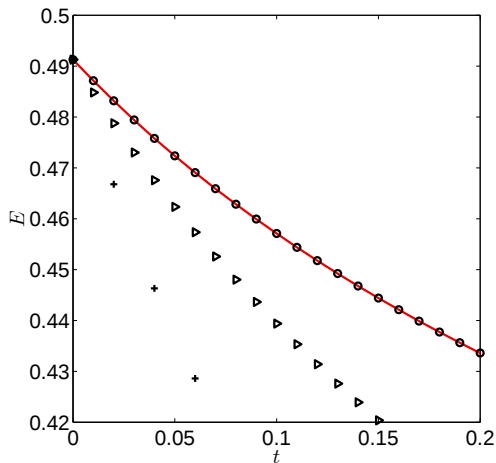


Numerical method:

- **Pseudo-spectral discretization:** full de-aliasing, parallelized
- **Helical-wave decomposition:** Coriolis diagonalization

Benefit of helical wave decomposition (HWD)

Decay of kinetic energy:



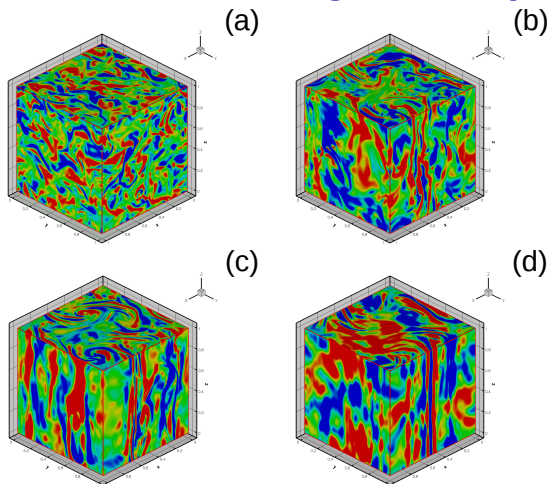
Rotation: $Ro = 10^{-2}$

RK4 with HWD (solid)
at $\Delta t = 3 \cdot 10^{-3}$

RK4 without HWD at
 $\Delta t = 7 \cdot 10^{-3}$ (+)
 $\Delta t = 5 \cdot 10^{-3}$ (▷)
 $\Delta t = 5 \cdot 10^{-4}$ (○)

More robust and larger time steps possible

Vertical structuring - vorticity

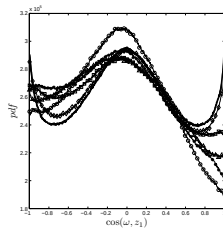


Decaying turbulence at $R_\lambda = 92$. Snapshot $t = 2.5$: $Ro = \infty$
(a), $Ro = 0.2$ (b), $Ro = 0.1$ (c), $Ro = 0.02$ (d)

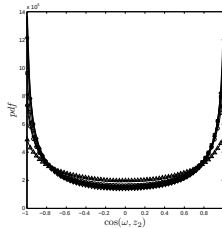
Geometrical statistics - alignment

Eigenvalues/vectors S_{ij} : $(\lambda_i, \mathbf{z}_i)$

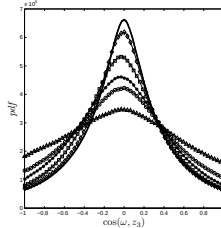
(positive) $\lambda_1 > \lambda_2 > \lambda_3$ (negative)



(1)



(2)



(3)

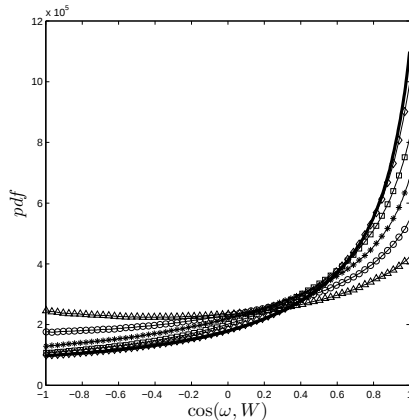
• $\cos(\omega, \mathbf{z}_i)$

$i = 1$ $Ro \rightarrow \infty$ – aligned or perpendicular, rotation reduces parallel alignment

$i = 2$ alignment ω and $\mathbf{z}_2$ – reduced by rotation

$i = 3$ $Ro \rightarrow \infty$ – mainly perpendicular, reduced preference with rotation

Geometrical statistics - vortex stretching



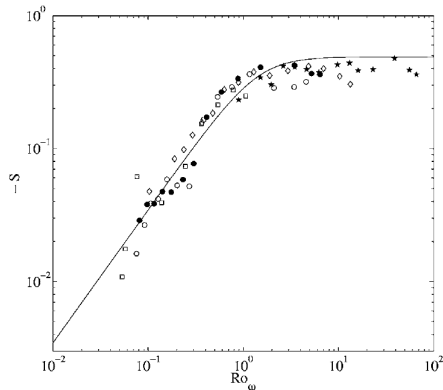
• $\cos(\omega, \mathbf{W})$

$Ro \rightarrow \infty$ vortex stretching dominant over vortex compression

$Ro \downarrow$ decreasing dominance vortex stretching

Depleted skewness at strong rotation

Experimental observation - reduced skewness at high $1/Ro$:

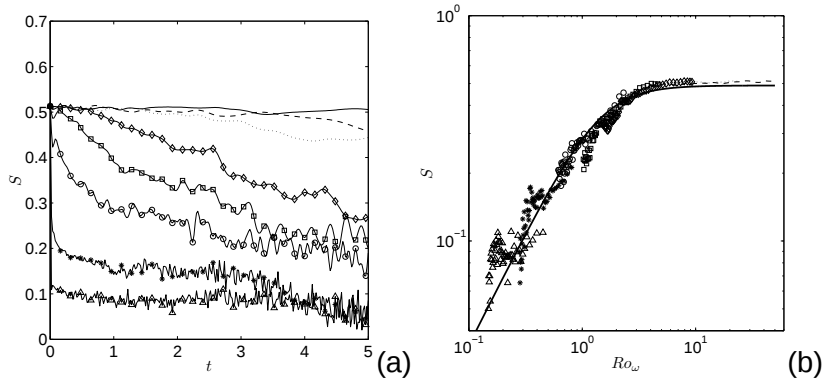


Cambon et al.

$$S = \frac{A}{(1 + 2Ro^{-2})^{1/2}} ; \quad S \sim Ro \sim \Omega^{-1} \text{ as } Ro \ll 1$$

Depleted skewness at strong rotation

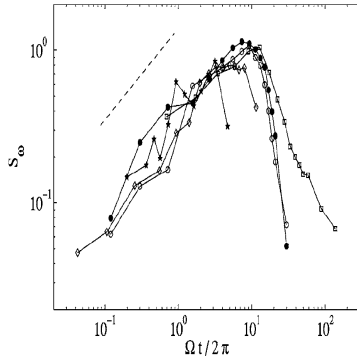
Simulated skewness: DNS at 512^3



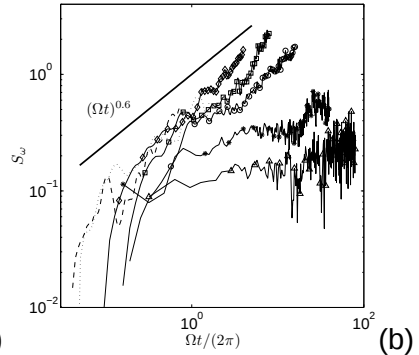
Velocity skewness versus t (a) and versus Ro (b)

General agreement Cambon and experiment

Vorticity skewness



(a)

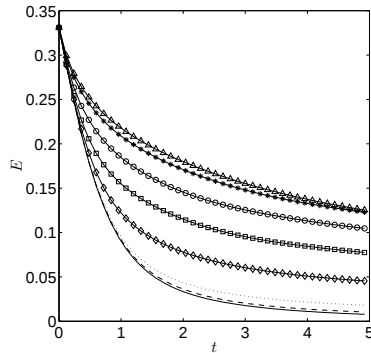


(b)

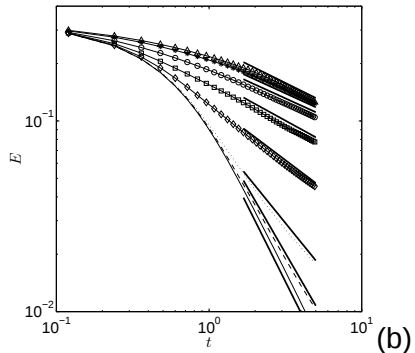
Experimental (a) and simulated (b) vorticity skewness at various rotation rates Ω

Algebraic decay of kinetic energy

DNS at 256^3 :



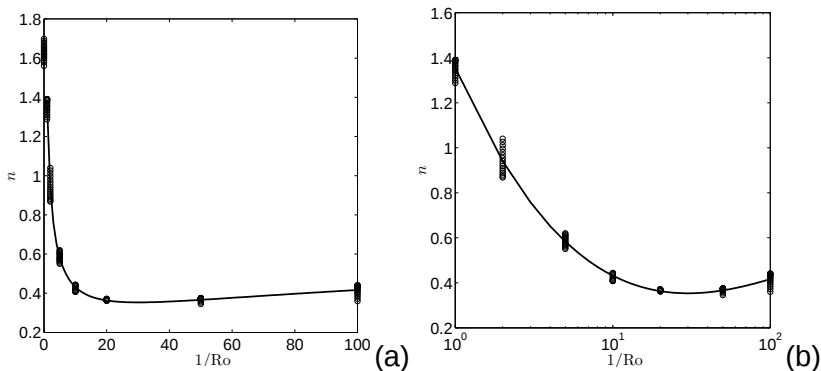
(a)



(b)

Decay of kinetic energy - slower decay at stronger rotation

Algebraic decay exponent

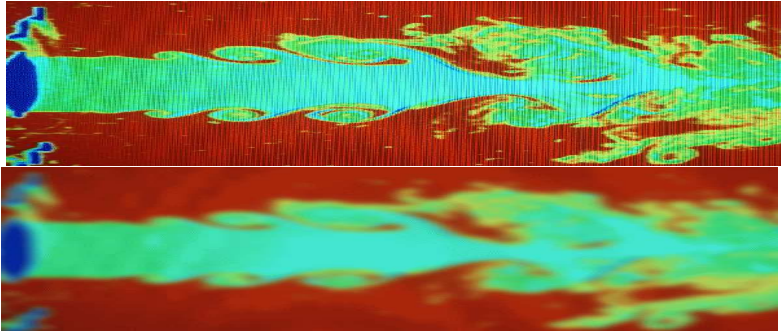


Exponent computed over various time-intervals $t_2 - t_1$.

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DNS and LES in a picture



Classical problem: wide dynamic range

- capture both large and small scales: resolution problem

$$N \sim Re^{9/4} ; W \sim Re^3$$

If $Re \rightarrow 10 \times Re$ then $N \rightarrow 175N$ and $W \rightarrow 1000W$

- Coarsening/mathematical modeling instead: LES

Filtering Navier-Stokes equations

Convolution-Filtering: filter-kernel G

$$\bar{u}_i = L(u_i) = \int G(x - \xi) u(\xi) d\xi$$

Application of filter:

$$\partial_t \bar{u}_i + \partial_j (\bar{u}_i \bar{u}_j) + \partial_i \bar{p} - \frac{1}{Re} \partial_{jj} \bar{u}_i - \frac{1}{Ro} (\bar{\mathbf{u}} \times \boldsymbol{\Omega})_i = -\partial_j \tau_{ij}$$

Turbulent stress tensor: $\tau_{ij} = \overline{u_i u_j} - \bar{u}_i \bar{u}_j$

Shorthand notation: Closure problem

$$NS(\mathbf{u}) = 0 \quad \Rightarrow \quad NS(\bar{\mathbf{u}}) = -\nabla \cdot \boldsymbol{\tau}(\mathbf{u}, \bar{\mathbf{u}}) \Leftarrow -\nabla \cdot \mathbf{M}(\bar{\mathbf{u}})$$

Closed LES formulation

$$\text{Find } v : \quad NS(v) = -\nabla \cdot \mathbf{M}(v)$$

Various closure models – (dynamic) eddy-viscosity, similarity, ...

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Regularization models

Alternative modeling approach

- 'first principles' – derive implied subgrid model
- achieve unique coupling to filter, e.g., Leray, LANS- α

Example: Leray regularization - alteration convective fluxes

$$\partial_t u_i + \bar{u}_j \partial_j u_i + \partial_i p - \frac{1}{Re} \Delta u_i = 0$$

LES template:

$$\partial_t \bar{u}_i + \partial_j (\bar{u}_j \bar{u}_i) + \partial_i \bar{p} - \frac{1}{Re} \Delta \bar{u}_i = -\partial_j (m_{ij}^L)$$

Implied Leray model

$$m_{ij}^L = L \left(\bar{u}_j L^{-1}(\bar{u}_i) \right) - \bar{u}_j \bar{u}_i = \overline{\bar{u}_j u_i} - \bar{u}_j \bar{u}_i$$

- Rigorous analysis available ($\sim 1930s$)
- Provides accurate LES model (~ 03)

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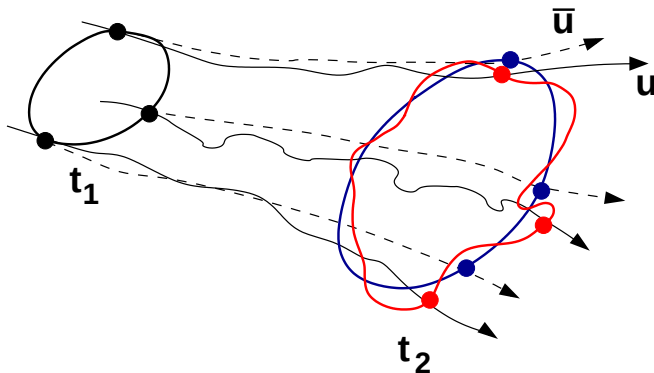
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NS- α regularization

Kelvin's circulation theorem

$$\frac{d}{dt} \left(\oint_{\Gamma(\mathbf{u})} u_j dx_j \right) - \frac{1}{Re} \oint_{\Gamma(\mathbf{u})} \partial_{kk} u_j dx_j = 0 \Rightarrow \text{NS} - \text{eqs}$$



Filtered Kelvin theorem

NS- α regularization

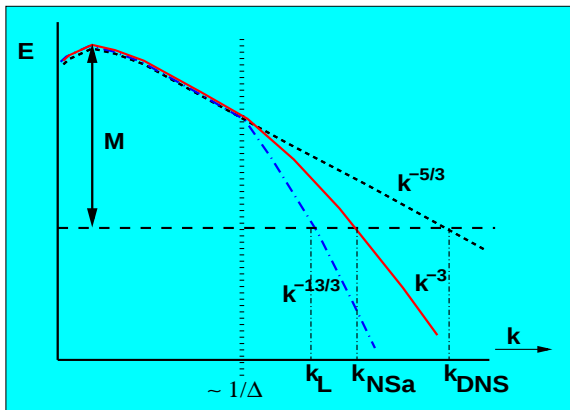
α -principle: Fluid loop: $\Gamma(\mathbf{u}) \rightarrow \Gamma(\bar{\mathbf{u}}) \Rightarrow$ Euler-Poincaré

$$\partial_t u_j + \bar{u}_k \partial_k u_j + u_k \partial_j \bar{u}_k + \partial_j p - \partial_j \left(\frac{1}{2} \bar{u}_k u_k \right) - \frac{1}{Re} \partial_{kk} u_j = 0$$

Rewrite into LES template: Implied subgrid model

$$\begin{aligned} \partial_t \bar{u}_i &+ \partial_j (\bar{u}_j \bar{u}_i) + \partial_i \bar{p} - \frac{1}{Re} \partial_{jj} \bar{u}_i \\ &= -\partial_j \left(\overline{\bar{u}_j u_i} - \bar{u}_j \bar{u}_i \right) - \frac{1}{2} \left(\overline{u_j \partial_i \bar{u}_j} - \bar{u}_j \partial_i u_j \right) = -\partial_j m_{ij}^\alpha \end{aligned}$$

Cascade-dynamics – computability



- $NS-\alpha$, Leray are dispersive
- Regularization alters spectrum – controllable cross-over as $k \sim 1/\Delta$: steeper than $-5/3$

Confined and nonconfined decay

Decay exponent n and spectral exponent p

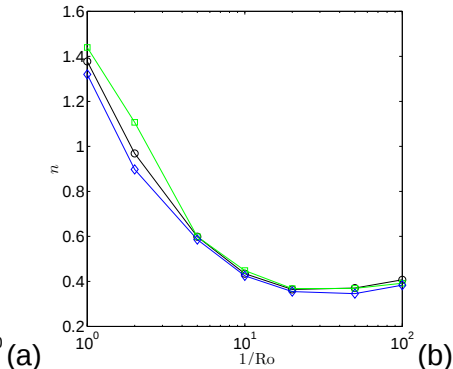
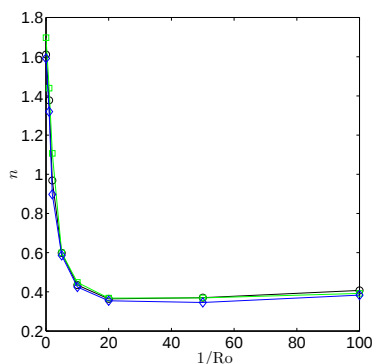
		Nonconfined	Confined
Nonrotating	$p = 5/3$	$n = 6/5$	$n = 2$
Rotating(I)	$p = 2$	$n = 3/5$	$n = 1$
Rotating(II)	$p = 3$	$n = 0$	$n = 0$

Rotating I – energy transfers time scale Ω^{-1}

Rotating II – totally inhibited energy transfer (Kraichnan)

LES: Algebraic decay at $R_\lambda = 92$

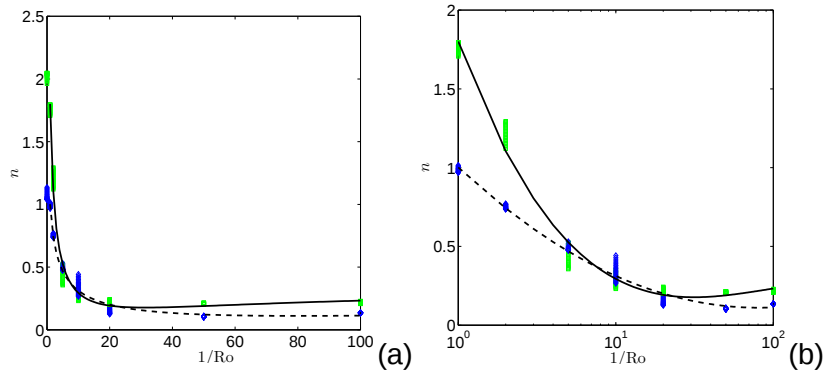
DNS at 256^3 – LES at 64^3 :



- Decay exponent well captured: DNS, Leray, LANS- α
- Evidence of nonconfined decay; Rotating I or II ?

LES: Algebraic decay at $R_\lambda = 200$

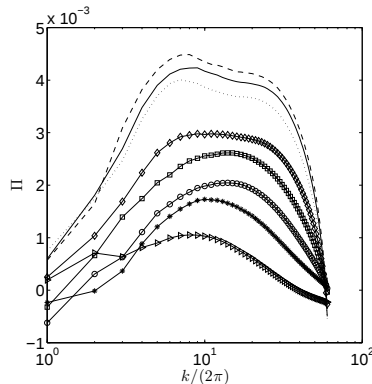
LES at 128^3 :



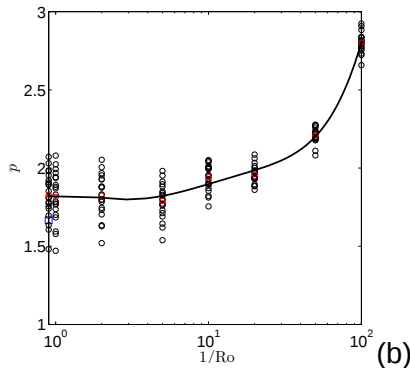
- Leray - dashed, LANS- α – solid
- Leray closer to ‘nonconfined’, LANS- α closer to ‘confined’
- Evidence Rotating II - Kraichnan

LES: Leray spectra at $R_\lambda = 200$

LES at 128^3 : totally inhibited energy transfer regime



(a)



(b)

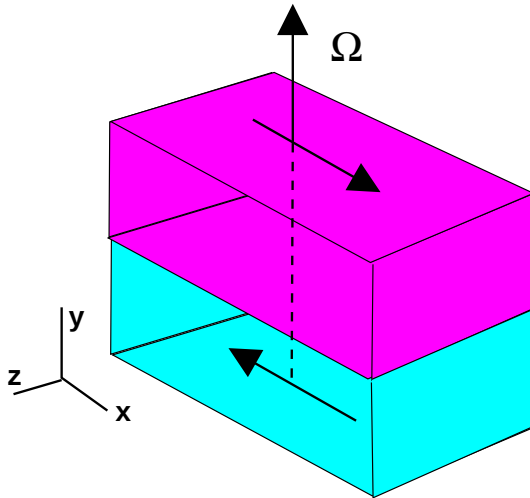
(a): Average transport-power spectra

(b): Energy-spectrum exponent showing transition
 $-5/3$ ($Ro \rightarrow \infty$) to -3 ($Ro \rightarrow 0$) – Kraichnan

Outline

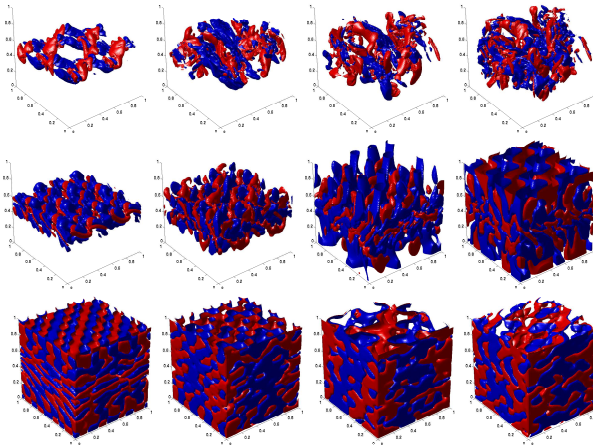
- 1 Rotation modulates turbulence
- 2 Rotation of incompressible fluid
- 3 Rotation at high Reynolds numbers
- 4 Rotation induced compressibility effects**
- 5 Concluding remarks

Compressible rotating mixing layer



- Consider range of Ro

Vorticity along axis of rotation: $\omega_2 = \partial_3 u_1 - \partial_1 u_3$



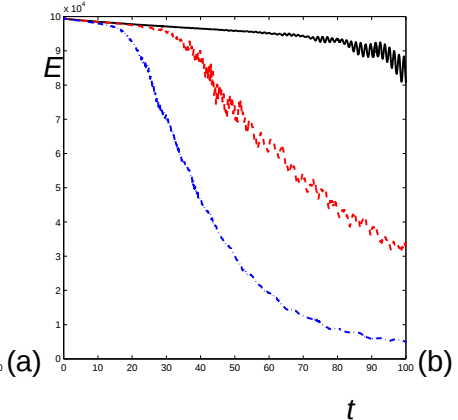
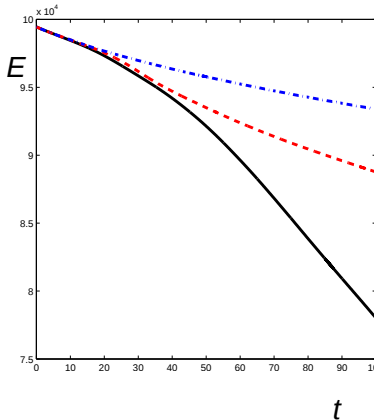
(30)

(60)

(80)

(90)

Kinetic energy decay



(a): $Ro = \infty$ (solid), $Ro = 10$ (dashed), $Ro = 5$ (dash-dotted)

(b): $Ro = 1$ (solid), $Ro = 0.5$ (dashed) and $Ro = 0.2$ (dash-dotted)

Contributions to decay rate

Total decay:

$$\frac{dE}{dt} = D - W \quad ; \quad D = \langle p \nabla \cdot \mathbf{u} \rangle \quad ; \quad W = \frac{1}{2Re} \langle \mathbf{S} : \mathbf{S} \rangle$$

D : pressure/velocity divergence, W dissipation

- incompressible: $D = 0$, $W \geq 0 \rightarrow$ monotonous decay of E
- compressible: $W \geq 0$, $D \neq 0 \rightarrow$ non-monotonous decay

Effect of rotation:

- no direct influence
- contributions to dD/dt and dW/dt

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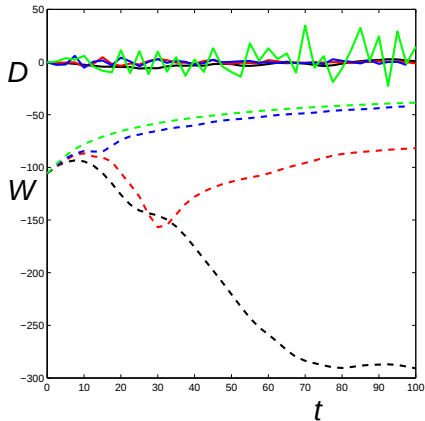
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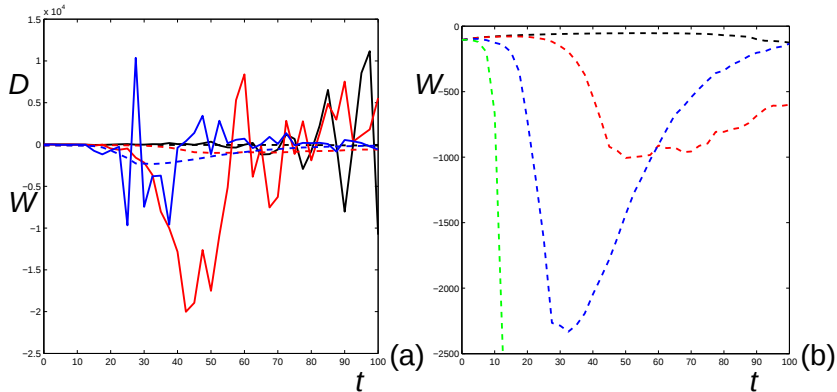
Slow rotation – incompressible limit



Dissipation and pressure/velocity divergence:

- $Ro = \infty$, $Ro = 10$, $Ro = 5$, $Ro = 2$
- reduced decay-rate as $Ro \downarrow$

Rapid rotation – compressibility



Dissipation W and pressure/velocity divergence D :

- $Ro = 1$, $Ro = 0.5$, $Ro = 0.2$, $Ro = 0.1$
- dominance D over W as $Ro \downarrow$
- non-monotonous decay

Outline

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Concluding remarks

- Established general correspondence DNS and experiments MMR
- Illustrated relevance helical wave decomposition
- Accurate predictions – Leray and LANS- α at $R_\lambda = 92$
- Leray at $R_\lambda = 200$ – exponent spectrum from $-5/3$ to -3 as $Ro \downarrow$: 3D - 2D
- Compressibility effects induced by Coriolis force: non-uniform decay, indirect rotation effect