

roots $\rightarrow$ positive roots, simple roots α_i , Fundamental weights

[see later] [12]

$$A_{ij} = \frac{2(\alpha_i, \alpha_j)}{(\alpha_j, \alpha_j)}, \quad 2 \frac{(\Lambda_i, \alpha_j)}{(\alpha_j, \alpha_j)} = \delta_{ij} \quad \Lambda = (A^{-1})_{ij} \alpha_j \quad \begin{matrix} \nearrow \text{[give rise to basic} \\ \text{rep of group]} \end{matrix}$$

The highest weight of irrep

[d] [see below]

$$\Lambda = \sum_{i=1}^K n_i \Lambda_i \quad \text{if } \Lambda = \Lambda_i \quad \text{each of its basic irrep}$$

[see later] [8]

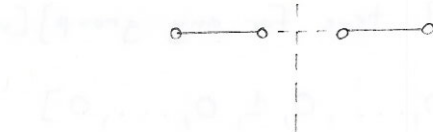
An: $\dim [0, \dots, 0, \underset{\substack{\uparrow \\ \text{K-th}}}{1}, 0, \dots, 0] = \binom{n+1}{K}$ K-th rank antisymmetric representation of $SU(n+1)$

$$\binom{n+1}{K} = \binom{n+1}{n+1-K}$$

[Prop binomial
coeffⁿ have
symmetry]

Set $N = n+1$

$$\dim [n_1, \dots, n_{N-1}] = \dim [n_{N-1}, \dots, n_1]$$



Example

$$A_3 = D_3 = SU(4), SO(6)$$

$$\dim [n_1, n_2, n_3] = \dim [n_3, n_2, n_1]$$

[antisym product]

$$\dim [1, 0, 0] = 4 = \dim [0, 0, 1] \quad \dim [0, 1, 0] = 6$$

[0, 1, 0]
//

$[n_1, n_2, n_3]$ is in the tensor product of $\text{Sym}^{n_1} [1, 0, 0] \otimes \text{Sym}^{n_2} (\Lambda^2 [1, 0, 0])$
 $\otimes \text{Sym}^{n_3} (\Lambda^3 [1, 0, 0])$

[0, 0, 1]

[When take K-th asym product $\rightarrow$
move 1 along rows of $[0, 0, 0s]$]

For a representation of $SU(N)$ of
Dynkin labels

[just deal with
algebra not group
- so only δ_{ij} , not
 η_{ij} at [d]]

$[n_1, \dots, n_{N-1}]$ the complex conj rep
is $[n_{N-1}, \dots, n_1]$

$$\left(\begin{matrix} m_1 & m_2 & m_3 & m_4 & m_5 \\ \square & \square & \square & \square & \square \\ \square & \square & \square & \square & \square \\ \square & \square & \square & \square & \square \\ \square & \square & \square & \square & \square \end{matrix} \right)^c \leftarrow \text{[complex conjugate]}$$

For $SU(4)$

$$(m_4=1) \quad (N-m_4=3)$$

$$\left(\begin{matrix} \square \\ \square \\ \square \end{matrix} \right)^c = \begin{matrix} \square \\ \square \\ \square \end{matrix}$$

$$(m_3=2) \quad (N-m_3=2)$$

$$\left(\begin{matrix} \square \\ \square \end{matrix} \right)^c = \begin{matrix} \square \\ \square \end{matrix}$$

$$\left(\begin{matrix} N-m_5 & N-m_4 & N-m_3 & \dots \\ \downarrow & \downarrow & \downarrow & \dots \\ N-m_2 & N-m_1 & N-m_0 & \dots \\ \square & \square & \square & \dots \\ \square & \square & \square & \dots \\ \square & \square & \square & \dots \\ \square & \square & \square & \dots \\ \square & \square & \square & \dots \end{matrix} \right) \quad \text{[corrected]}$$

[Originally
 $\dots m_1, \dots m_2, \text{etc.}$
changed to
 $\dots m_5, \dots m_4, \text{etc.}$]

All irreps $[0, n, 0]$ are real, symmetric traceless with rank tensors of $SO(6)$

$$\begin{matrix} (2,1) \\ \left(\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} \right)^c = \begin{matrix} (4-1, 4-2) \\ = (3,2) \\ \begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} \end{matrix}$$

For $SU(6)$ $[0, 0, n, 0, 0]$ is real
i.e. $[1, 1, 0]^c = [0, 1, 1]$

$[n_1, n_2, n_3 \dots n_3, n_2, n_1]$ in general

[All asym reps are real]

[All this A_n]

[Fund weights correspond to K -th rank asym rep]

$[A] \rightarrow [X]$ true for any group [now apply to B_n]

B_n $[0, \dots, 0, 1, 0, \dots, 0]$ $\circ \dots \circ \Rightarrow \circ$

$\dim [0, \dots, 0, 1] = 2^n$ spinor of $SO(2n+1)$

Examples

$B_1 = C_1 = A_1 \cong SU(2)$ $2^1 = 2$ fund. rep

$\det(A_{B_n}) = 2$ $SO(3)$
 $Sp(1)$

[Stern-Gerlach experiment: spin relate to ang. momenta + Fermi stats]

2 types of irreps $[n_1, \dots, n_n]$, if n_n is even bosonic
odd fermionic

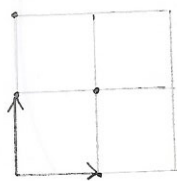
tensor product of

bosonic $\times$ bosonic gives bosonic
" ferm ferm ferm
ferm $\times$ ferm bosonic

[stay in bosonic lattice]

[stay in fermionic lattice]

[move from ferm to bosonic]



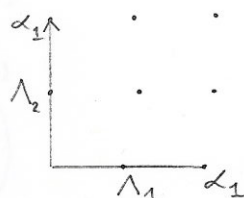
[can use determinant to find area]
[see note 1, p.4]

A_1 : $\Lambda = \frac{1}{2} \alpha$
[Fund weight] [true root]

[because of properties of determinant could find solution]



$D_2 = A_1 \times A_1$



[Area 4 times as large for D_2 as area picked out by roots]

$$[0, \dots, 0, 1, 0, \dots, 0]$$

K -th rank antisymmetric rep. of $SO(2n+1)$
of dimension $\binom{2n+1}{K}$ real rep

[vector space
always real
by construction]

↑ [referring back to {bottom} root diagrams] {p2}

$[1, 0]$ spinor

$[0, 1]$ "

$[2, 0]$ self dual
2nd rank
antisym
rep of $SO(4)$

a representation that appears in physics $[1, 0, \dots, 0, 1]$

↑
vector

↑
spinor

~ gravitino

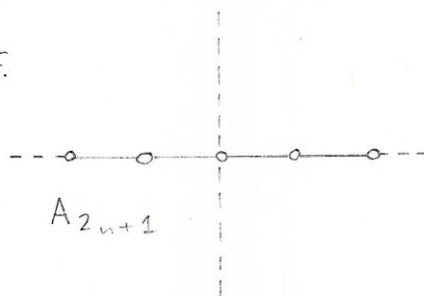
[appears
in $d \geq 4$
dim-context
supgrav]

dim

$$\left. \begin{array}{l} B_2: [0, 1] \quad 4, \text{pseudoreal} \\ B_3: [0, 0, 1] \quad 8, \text{real} \\ B_4: [0, 0, 0, 1] \quad 16, \text{real} \end{array} \right\} \text{spinor reps}$$

$$C_n \circ \cdots \circ \circ \leftarrow \circ$$

cf.



$$A_5 \simeq SU(6) \quad [1, 0, 0, 0, 0]$$

$$\chi_{[1, 0, 0, 0, 0]} = \gamma_1 + \frac{\gamma_2}{\gamma_1} + \frac{\gamma_3}{\gamma_2} + \frac{\gamma_4}{\gamma_3} + \frac{\gamma_5}{\gamma_4} + \frac{1}{\gamma_5}$$

set the fugacities of C_3 to be Z_i

$$Z_1 = \gamma_1 = \gamma_5$$

$$Z_2 = \gamma_2 = \gamma_4$$

$$Z_3 = \gamma_3$$

$$\chi_{[1, 0, 0]_{C_3}} = Z_1 + \frac{Z_2}{Z_1} + \frac{Z_3}{Z_2} + \frac{Z_2}{Z_3} + \frac{Z_1}{Z_2} + \frac{1}{Z_1} \quad \text{pseudoreal}$$

[singlet
derived
from
sym or asym
-determine
pseudoreal
or not]

$$[0, \dots, 0, 1, 0, \dots, 0]_{C_n}$$

K -th

traceless K -th rank
antisymmetric rep of $Sp(n)$

[How know traceless?]

[not simple dim
to write down
-not binomial
since traceless]

Recall for $SO(n)$ 2 vectors $V_i, W_i \quad i=1, \dots, n$

$$V_i W_j + V_j W_i \quad \text{second rank symmetric}$$

$$V_i W_i \quad \text{is a singlet}$$

$Sp(n)$ V_α, W_β antisymmetric $V_\alpha W_\beta - V_\beta W_\alpha$ reducible

$$J^{\alpha\beta} V_\alpha W_\beta$$

[see traceless
by analogy]

Note 1 Relation Between Determinant

And Area (general formulation)

In 2D: area of a parallelogram spanned by

$$\underline{a} = a_1 \underline{i} + a_2 \underline{j}$$

$$\underline{b} = b_1 \underline{i} + b_2 \underline{j}$$

is given by

$$\|\underline{a} \times \underline{b}\| = \left| \det \begin{pmatrix} a_1 & a_2 \\ b_1 & b_2 \end{pmatrix} \right|$$

Physics: A_μ , $\mu=0\dots 3$ so we are in $SO(4)$
 i.e. $D_2 = A_1 \times A_1$

4D rep: $[3,0]$, $[0,3]$ or $[1,1]$, or $[1,0] + [0,1]$

spin $\frac{3}{2}$ rep for
 one $SU(2)$
 do not want this
 for general vectors

idem

the vector
 rep we want

2 spinors, which
 do not work
 properly as well

For the $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$

∂_μ transf in $[1,1]$ (i.e. as a vector A_μ)

$F_{\mu\nu}$ — " — " — $\Lambda^2[1,1]$ antisym product of two vectors

$$\dim \Lambda^2[1,1] = \binom{4}{2} = 6$$

Compute $\Lambda^2[1,1]$ in irrep with PE_F

$$\text{or } f(X_1, X_2) = (X_1 + X_1^{-1})(X_2 + X_2^{-1}) = \chi_{[1,1]}(X_1, X_2)$$

$$\text{then } \frac{1}{2} f^2(X_1, X_2) - \frac{1}{2} f(X_1^2, X_2^2) = (2,0)(0,2)(0,0)(0,0)(0,-2)(-2,0)$$

$$\text{from } (1,1)(1,-1)(-1,1)(-1,-1)$$

Highest weight: $[2,0]$

$$\text{so } \Lambda^2(1,1) = [2,0] + [0,2] \quad (\text{sym in the exchange of } [\cdot, \cdot])$$

So $F_{\mu\nu}$ is reducible

It's a reducible rep of $SO(4)$: the adjoint rep of $SO(4)$

Can then be written

$$F_{\mu\nu} = F_{\mu\nu}^1 + F_{\mu\nu}^2 \quad \text{which transforms without mixing}$$

$$\text{Using } \varepsilon_{\mu\nu\rho\sigma} F^{\rho\sigma} \equiv \tilde{F}_{\mu\nu}, \text{ we can write } F_{\mu\nu}^\pm \equiv \frac{1}{2}(F_{\mu\nu} \pm \tilde{F}_{\mu\nu})$$

$$\text{and the } F_{\mu\nu} = F_{\mu\nu}^+ + F_{\mu\nu}^-$$

Does this work?

$$\varepsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^+ = \varepsilon^{\mu\nu\rho\sigma} \left(\frac{1}{2} F_{\mu\nu} + \frac{1}{2} \varepsilon_{\mu\nu\alpha\beta} F^{\alpha\beta} \right) = \frac{1}{2} \tilde{F}^{\rho\sigma} + \frac{c}{2} \varepsilon^{\rho\sigma\alpha\beta} F_{\alpha\beta}$$

$$= F^{\rho\sigma}$$

shows
 that

i.e. $F_{\mu\nu}^+$ has an eigenvalue +1 under $\varepsilon^{\mu\nu\rho\sigma}$ as well as

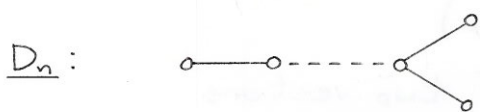
$$F_{\mu\nu}^- \quad \text{---} \quad \text{---} \quad \text{---} \quad \text{---} \quad -1 \quad \text{---} \quad \text{---}$$

We can identify these self dual and anti-self dual tensors $F_{\mu\nu}^+$ and $F_{\mu\nu}^-$ as the elements transforming in $[2,0]$ and $[0,2]$

The middle dimensional antisym tensor:

$\frac{2n}{2} = n$ $\swarrow$ $\searrow$ $SO(2n)$: the n -th rank antisym rep of $SO(2n)$ is reducible into SD and ASD tensors

$$\text{Self dual} \Leftrightarrow F^+ = 0, \text{ antiself dual} \Leftrightarrow F^- = 0$$



$$SO(2n): \text{ basic rep are } [0, \dots, \underset{\substack{\uparrow \\ k^{\text{th}}}}{1}, \dots, 0]$$

For $K=n-1$ or n : spinor rep of $SO(2n)$, dim 2^{n-1} distinguished by calling spinor and spinor (!)

Dirac spinor (on $SO(4)$) is reducible in $[1,0] + [0,1]$

(because massless part)

called two Weyl spinors, the L and R ones

(whereas Dirac spinor is L-R sym)

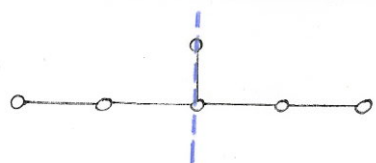
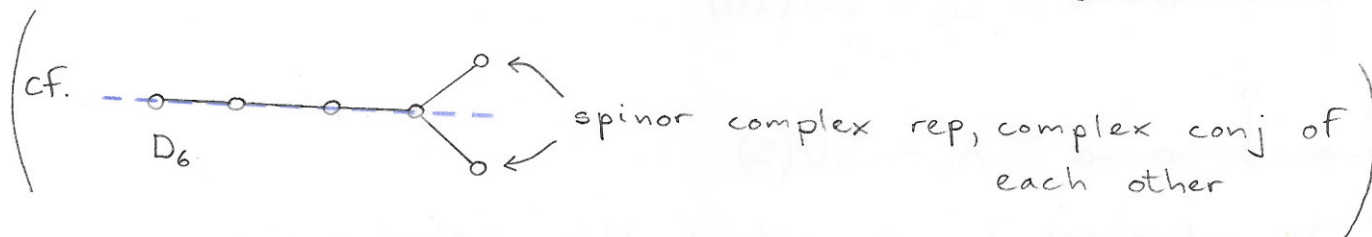
Let $A_{D_n} = 4$ so set the last two indices to be n_{n-1}, n_n being even or odd (each)

$$\text{giving } [n_1 \dots n_{n-1}, n_n] \otimes [m_1 \dots m_{n-1}, m_n] = [\dots n_{n-1} + m_{n-1}, n_n + m_n]$$

Selection rules:

$$[\dots \text{odd}, \text{odd}] \otimes [\dots \text{odd}, \text{odd}] = [\dots \text{even}, \text{even}]$$

odd, even	odd, odd	even, odd
$\vdots$	$\vdots$	$\vdots$

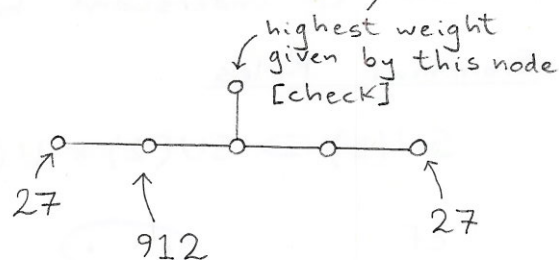
Decomposition of exceptional groups E_i ($i=6,7,8$) E_6 :diag symmetric under reflection,
so under complex conjugation

$$\dim [1, 0 \dots 0]_{E_6} = 27 = \dim [0 \dots 0, 1]_{E_6}$$

traceless antisym
[i.e. adjoint]

$$\dim \text{Adj}(E_6) = 78$$

$$\dim [0, 1, 0 \dots 0] = 912$$



E_7 : $\dim \text{Funda } E_7 = 56$ pseudoreal; $\dim \text{Adj } E_7 = 133$

E_8 : $\dim \text{Adj } E_8 = 248$

Recall: every dot in Dynkin diag is simple root and Funda weight
(one α_i corresponds to one Λ_j :

$$\Lambda_i = (A^{-1})_{ij} \alpha_j \quad \text{with} \quad A_{ij} = \frac{2(\alpha_i, \alpha_j)}{(\alpha_j, \alpha_j)} \quad \text{[i.e. roots allow us to form the Cartan matrix]}$$

$$\text{i.e.} \quad 2 \frac{(\Lambda_i, \alpha_j)}{(\alpha_j, \alpha_j)} = S_{ij}$$

each node is represented by n
and we can write:

$$\dim(n_1 \ n_2 \ n_3 \ n_4 \ n_5 \ n_6)$$

every rep is represented by a highest weight Λ

decomposition in Λ_i :

$$\Lambda = \sum_i n_i \Lambda_i, \quad \text{with } \Lambda_i \text{ the Funda weights}$$

Dynkin diagram contains all possible roots

GUT:

$$SU(5) \supset SU(3) \times SU(2) \times U(1)$$

$$SO(10)$$

$$E_6$$

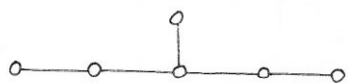
How to couple a scalar field to a gauge field?
Use covariant derivative

Gauge Field Coupling $R = \text{representation}$

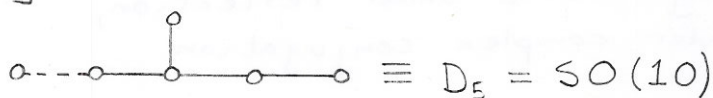
$$D_\mu \phi_R = \partial_\mu \phi_R + ig A_\mu \phi_R \quad \frac{1}{2} |D_\mu \phi_R|^2$$

ϕ is a vector space; A_μ transforming it, rep gives the interaction, i.e. the action of A_μ on ϕ , and the mass of fields

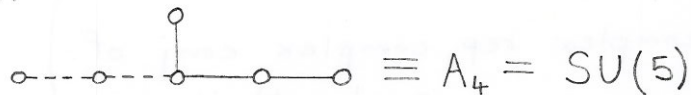
Indeed, consider Dynkin diagram for E_6



taking a subset



then

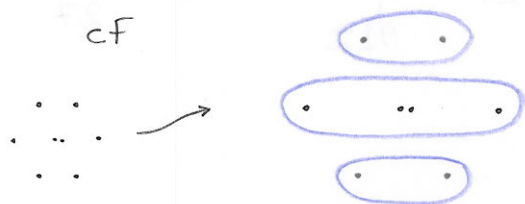


but need to understand how to extend the mechanisms: branching

Branching Rules

$$SU(3) \supset SU(2) \times U(1) \leftarrow \left\{ \begin{array}{l} \text{Suppose we have a group } SU(3) \\ \text{and we want a subgroup decomposition} \\ SU(2) \times U(1) \end{array} \right\}$$

cf



in $SU(3)$ rep

decomposed into $SU(2)$ different rep

or 3×3 matrix:

$$\left(\begin{array}{c|c} 2 \times 2 & 1 \times 2 \\ \hline 2 \times 1 & 1 \times 1 \end{array} \right)$$

Hence $3 \rightarrow 2 + 1$

but charged under the other gp:

$$2_1 + 1_{-2}$$

↑ ↑
charges

2nd rank sym of $SU(3) = 6 \rightarrow ?$

Use the fugacities:

$$\chi_{[1,0]}(Y_1, Y_2) = Y_1 + \frac{Y_2}{Y_1} + \frac{1}{Y_2} = \chi_3$$

$$3 = \begin{pmatrix} \cdot & \cdot \\ \cdot & \cdot \end{pmatrix} \equiv \left. \begin{array}{l} \text{doublet} \\ \text{singlet} \end{array} \right\} \text{with different charges}$$

$$\begin{cases} \chi_3 = q^1 \chi_2 + q^{-2} \chi_1 \\ 3 = 2_1 + 1_{-2} \end{cases}$$

Note: $SU(2) \left\{ \begin{array}{l} \text{Fugacity} \\ \text{highest weight} \end{array} \right\} = X, \quad U(1) \left\{ \begin{array}{l} \text{fugacity} \\ \text{charge} \end{array} \right\} = q,$

set $\cdot$ $SU(2)$ fugacity to X
 $\cdot$ $U(1)$ " " " " q

choosing this mapping: $\uparrow$
 $\chi_3 = q \left(\frac{1}{X} + X \right) + \frac{1}{q^2}$

giving:

$$Y_1 = qX, \quad \frac{Y_2}{Y_1} = \frac{q}{X}, \quad \frac{1}{Y_2} = \frac{1}{q^2} = \left(qX \cdot \frac{q}{X} \right)^{-1} = (q^2)^{-1} \quad \text{OK}$$

Fugacity map:

$$\boxed{Y_1 = qX, Y_2 = q^2} \quad \text{for the branching rules}$$

not choosing:

X and $\frac{1}{X}$ symmetry for $SU(2)$ properties ($qX \Rightarrow \frac{q}{X}$)

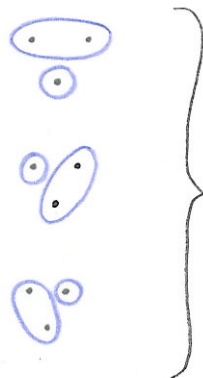
product of terms with X or $\frac{1}{X} =$ terms without it $\Rightarrow \frac{1}{q^2}$

Fugacity map not unique: 3 possible maps:

$$(1) \quad Y_1 = qX, \quad Y_2 = q^2$$

$$(2) \quad Y_1 = qX, \quad Y_2 = \frac{X}{q}$$

$$(3) \quad Y_1 = \frac{q}{X}, \quad Y_2 = q^2$$



Any of these will solve

$$Y_1 + \frac{Y_2}{Y_1} + \frac{1}{Y_2}$$

$$= q \left(X + \frac{1}{X} \right) + \frac{1}{q^2}$$

For 6:

• • • triplet
• • = doublet
• singlet

and $\chi_{[2,0]} = Y_1^2 + \frac{Y_2^2}{Y_1^2} + \frac{1}{Y_2^2} + Y_2 + \frac{Y_1}{Y_2} + \frac{1}{Y_1}$

↑
Dynkin labels
for E_6

with the first map: ($Y_1 = qX, Y_2 = q^2$)

$$\begin{aligned} \chi_{[2,0]} &= q^2 X^2 + \frac{q^4}{q^2 X^2} + \frac{1}{q^4} + q^2 + \frac{qX}{q^2} + \frac{1}{qX} \\ &= q^2 \left(X^2 + 1 + \frac{1}{X^2} \right) + \frac{1}{q} \left(X + \frac{1}{X} \right) + \frac{1}{q^4} \end{aligned}$$

$$\text{i.e. } \chi_{[2,0]} = q^2 \chi_{[2]} + q^{-1} \chi_{[1]} + q^{-4} \chi_{[0]}$$

$$\Rightarrow 6 = 3_2 + 2_{-1} + 1_{-4}$$

Once you have chosen a map there are 3 ways of implementing this map (in this case).

For χ_n there are n ways of choosing this map.

Branching rules

$$\left. \begin{array}{l} SU(3) \supset SU(2) \times U(1) \\ 3 \rightarrow 2_1 + 1_{-2} \end{array} \right\} \text{recap of lecture 28}$$

Fugacity map

e) Now: $E_6 \supset SO(10) \supset SU(5) \supset SU(3) \times SU(2) \times U(1)$

[Used to be popular
For SM representⁿ]

Decompose 5D vector:
→ 3 + 2D vector
(i.e. decomp of
funda of $SU(5)$)

$$5 \rightarrow (3, 1)_2 + (1, 2)_{-3}$$

triplet of $SU(3)$ = singlet under $SU(2)$ (a)
doublet under $SU(2)$ = singlet under $SU(3)$ (b)

$$\left(\begin{array}{c} \vdots \\ \vdots \\ \vdots \end{array} \right) \left\{ \begin{array}{l} 3\text{-d} \\ 2\text{-d} \end{array} \right.$$

(b) [singlet $SU(3)$
doublet $SU(2)$]

Cartan generators

$$O = \text{Tr} \left(\begin{array}{c|c} a & \\ \hline & a \\ \hline & & b \\ \hline & & & b \end{array} \right)$$

chosen to be proportional to

$$\left(\begin{array}{c|c} I_{3 \times 3} & 0 \\ \hline 0 & I_{2 \times 2} \end{array} \right)$$

generator $\Rightarrow$ traceless

$$3a + 2b = 0$$

$$a = 2, b = -3$$

[since normalizⁿ
is a choice]

[choose these]
[then write
2, -3 @
(a) & (b)]

[Cartan gen's -
trace zero]

[singlet carries
 $U(1)$ charges
- arbitrary
normalization
- only care e-
charge some
integer # times
some value
 $\Rightarrow$ deal w lattice
for non-Abelian
groups, though
this is Abelian]

[Fugacity map tells decomp for any repⁿ]

Characters $\chi_{[1,0,0,0]_{SU(5)}} = \chi_{[1,0]_{SU(3)}} \chi_{[0]_{SU(2)}} q^2 + \chi_{[0,0]_{SU(3)}} \chi_{[1]_{SU(2)}} q^{-3}$

$$[1,0,0,0]_{SU(5)} = \frac{Y_1}{Y_1} + \frac{Y_2}{Y_1} + \frac{Y_3}{Y_2} + \frac{Y_4}{Y_3} + \frac{1}{Y_4}$$

Note: using
[$n_1 \dots n_n$] For
 $\chi_{[n_1 \dots n_n]}$

[decomp want equal to this repⁿ - intro fugacities]

$$= q^2 \left(Z_1 + \frac{Z_2}{Z_1} + \frac{1}{Z_2} \right) + q^{-3} \left(X + \frac{1}{X} \right)$$

$SU(3)$ fugacities Z_1, Z_2

$SU(2)$ " X

$U(1)$ " q

Pick a fugacity rep:

$$Y_1 = q^2 Z_1 \quad (\text{choice}) \quad (1^{\text{st}} \text{ monomial}) \quad [\text{funda of } SU(5) = \text{funda of } SU(3) \times q^2 (\text{charge})]$$

$$\Rightarrow Y_2 = Y_1 \cdot \frac{q^2}{Z_1} Z_2 = q^4 Z_2 \quad (2^{\text{nd}} \text{ monomial}) \quad [\text{do fugacity map then find interpretⁿ}]$$

[character
fundamental
 $\times$ character
singlet]

$$Y_3 = Y_2 \cdot \frac{q^2}{Z_2} = q^6 \quad (3^{\text{rd}} \text{ monomial})$$

$$Y_4 = q^3 X \quad (5^{\text{th}} \text{ monomial})$$

[monomial]
[ternomial]

[4th rank antisym $SU(5)$ = map to $SU(2)$]

Check: $\frac{Y_4}{Y_3} = \frac{q^3 X}{q^6} = \frac{X}{q^3} \quad (4^{\text{th}} \text{ monomial}) \quad [\text{i.e. OK}]$

$[Y_1: SU(5) - \text{funda}]$

$[Y_2: SU(5) - 2^{\text{nd}} \text{ rank anti-funda}]$

$[Z_2: SU(3) - 2^{\text{nd}} \text{ rank anti-funda}]$

$[Y_3 - 3^{\text{rd}} \text{ rank anti-sym rep } SU(5)]$

$\Rightarrow 3 \times \text{charge } U(1) \text{ charge}]$

$[3^{\text{rd}} \text{ rank rep} \rightarrow \text{singlet}]$

$$\bar{5} \rightarrow (\bar{3}, 1)_{-2} + (1, 2)_3 \quad (\text{because } \bar{2}=2)$$

[cf γ_4 -possible to decompose to $SU(2)$ repⁿ] [i.e. $X_{q^3} = \gamma_4$ as above $\Rightarrow$ coherent]

[can tell how $SU(5)$ decompose to reps of $SU(3)$]

[4d lattice $SU(5)$ -different ways to choose sublattices also choose fugacities $\rightarrow$ natural basis, γ_i correspond to fund. weight etc. decomp doesn't depend on fugacity maps]

$$10: [0, 1, 0, 0]_{SU(5)} = \left(\gamma_2 + \frac{\gamma_3 \gamma_1}{\gamma_2} + \frac{\gamma_4 \gamma_1}{\gamma_3} + \frac{\gamma_1}{\gamma_4} \right)$$

[2nd rank ASym rep $SU(5)$]
[" " " & Sym simple - add squares on the sym]

$$+ \left(\frac{\gamma_3}{\gamma_1} + \frac{\gamma_4 \gamma_2}{\gamma_3 \gamma_1} + \frac{\gamma_2}{\gamma_1 \gamma_4} \right)$$

[10 monomials]

[using $\gamma_3^{-1} = q^{-6}$]

$$\underbrace{\gamma_2}_{q^4 \bar{z}_2 (c)} + \underbrace{\frac{\gamma_3 \gamma_1}{\gamma_2}}_{q^4 \frac{z_1}{z_2} (d)} + \underbrace{\frac{\gamma_4 \gamma_2}{\gamma_3 \gamma_1}}_{q^3 \times \frac{q^2}{q^6} z_1} + \dots + \underbrace{\frac{\gamma_1}{\gamma_4}}_{q^{-6}}$$

(By taking AS terms of product $\chi_{[1,0,0,0]} \chi_{[1,0,0,0]}$)

[dim = 3 + (3x2) + 1 = 10]

$$10 \rightarrow (\bar{3}, 1)_4 + (3, 2)_{-1} + (1, 1)_{-6}$$

[just by writing down 1st term] i.e. (c)

$\frac{1}{q} \times z_1$ [this gives the $(3, 2)_{-1}$ term]

[(d) part of anti-fund rep - only need pick some to get repⁿ]

Interpretation

$[(3, 2)_{-1}]$: Fit naturally with LH quarks

$[(\bar{3}, 1)_4]$: " " " " " RH quarks

$[(1, 1)_{-6}]$: take to predict massive particle not yet found [nature not LR symmetric]
[or RH electron $\rightarrow$ singlet under both]

[adjoint $SU(5)$ contains adjoint $SU(3)$ and $SU(2)$]

[leptons - singlet under $SU(3)$ and doublet under $SU(2)$]

trace:

3 weights	with $q=4 \Rightarrow +12$	} traceless
3x2 " "	" " $q=-1 \Rightarrow -6$	
1 " "	" " $q=-6 \Rightarrow -6$	

We have:

$$U(N) \supset U(n_1) \times U(n_2) \quad n_1 + n_2 = N$$

$$SO(N) \supset SO(n_1) \times SO(n_2)$$

$$Sp(N) \supset Sp(n_1) \times Sp(n_2)$$

[Rank:] $N-1$

$$SU(N) \supset S(U(n_1) \times U(n_2)) \cong \overbrace{SU(n_1)}^{n_1-1} \times \overbrace{SU(n_2)}^{n_2-1} \times \overbrace{U(1)}^1$$

[Removing Cartan element prop to id (non-zero trace) on both sides]

[natural embedding
U - complex vector space
SO - real vector space
Sp - 2n dim real vector space]

Note: can go to larger groups
 $SO(10) \supset SU(5)$ } see (e)
 $E_6 \supset SO(10) \dots$ } above

[removing traceless subgroup in Cartan algebra]

Branching Rules:

$G \supset H$ and write $R_G = \sum_i R_{H,i}$ with fugacity maps

We have:

$$U(N) \supset \prod_i U(n_i), \text{ with } \sum_i n_i = N$$

$$SU(N) \supset U(1)^{K-1} \prod_{i=1}^K SU(n_i), \text{ idem}$$

$$SO(N) \supset \prod_i SO(n_i), \text{ idem}$$

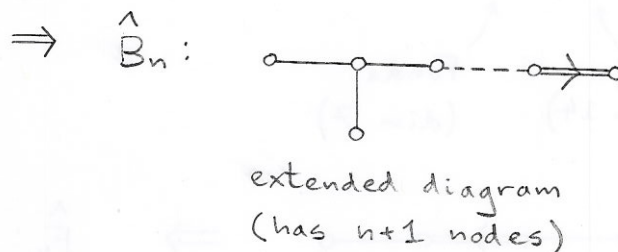
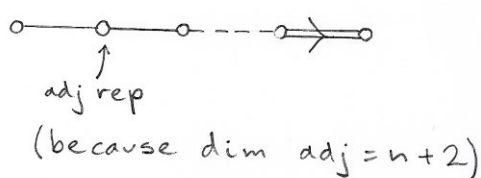
$$Sp(N) \supset \prod_i Sp(n_i), \text{ idem}$$

Dynkin method: to compute the decomposition (for classical and exceptional groups) use extended Dynkin diagram (for affine Lie algebra):

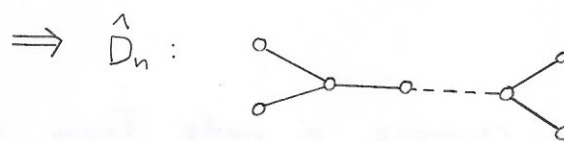
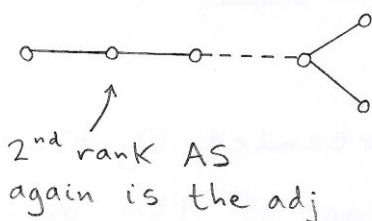
add one node that has a scalar product with the node of the adjoint rep

Examples

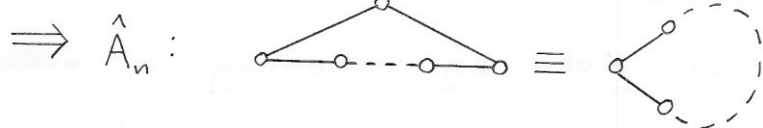
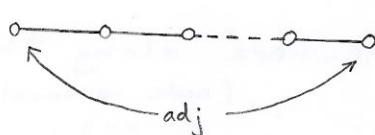
B_n :



D_n :

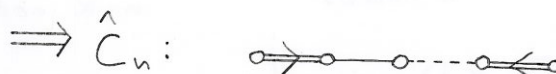
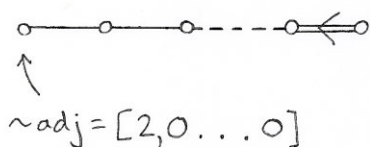


A_n :



(more symmetrically)

C_n :

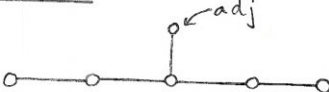
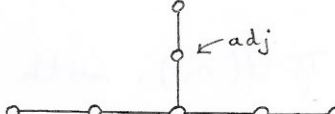


because for $SO(n)$, $\text{adj} \equiv 2^{\text{nd}}$ rank AS

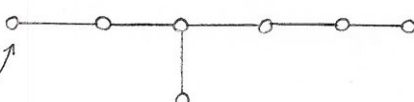
For $Sp(n)$, $\text{adj} \equiv 2^{\text{nd}}$ rank Sym

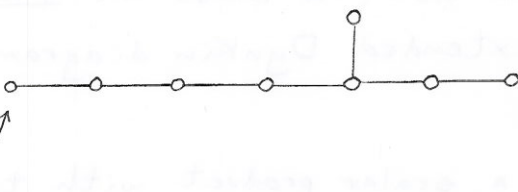
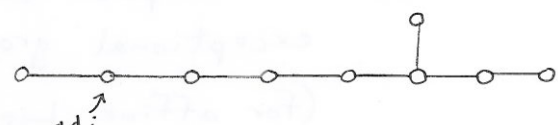
and because it is $[2, 0 \dots 0]$, we will put 2 lines from the node corresponding to $[1, 0 \dots 0]$

Exceptionals

E_6 :  $\Rightarrow$ $\hat{E}_6$: 

[notes copied from board p.5 $\rightarrow$ begin here]

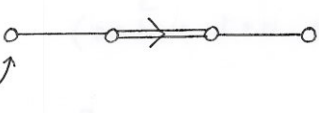
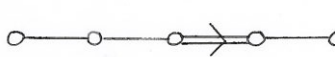
E_7 :  $\Rightarrow$ $\hat{E}_7$: 

E_8 :  $\Rightarrow$ $\hat{E}_8$: 

(where is the adjoint? can compute the dim)

G_2 :  $\Rightarrow$ $\hat{G}_2$: 

adj (dim 14) Funda (dim 7)

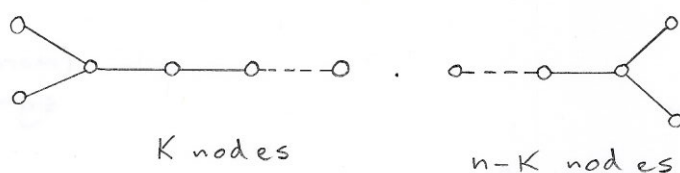
F_4 :  $\Rightarrow$ $\hat{F}_4$: 

adj

Prescription: remove a node from the extended Dynkin diag $\hat{G}$
the remaining diag is a subgp of the original gp G

Example

$\hat{D}_n$ gives by removing a node somewhere along the line
[node removal indicated by "..."]



$$\Rightarrow D_n \supset D_K \times D_{n-K}$$

$$\text{i.e. } SO(2n) \supset SO(2K) \times SO(2n-2K)$$

$$\begin{array}{c} \circ \\ \circ \end{array} \cdot \begin{array}{c} \circ \quad \circ \quad \circ \\ | \quad | \quad / \backslash \\ \circ \quad \circ \quad \circ \end{array} \Rightarrow D_n \supset SO(2n-4) \times \underbrace{SU(2) \times SU(2)}_{SO(4)}$$

Note that this method is not complete and useful mainly for exceptional cases

cf $\hat{A}_n$: gives A_n only $\Rightarrow$ useless

• $\hat{B}_n$ gives  = D_K and B_{n-K}

i.e. $SO(2n+1) \supset SO(2k) \times SO(2n-2k+1)$

For the last node, gives $SO(2n+1) \supset SO(2n)$

- $\hat{G}_2$ gives G_2 the trivial answer (via $\circ \Rightarrow \circ$)
- $A_1 \times A_1 = SU(2) \times SU(2)$ (via $\circ \cdot \circ$)
- $A_2 = SU(3)$ (via $\circ - \circ \cdot$)

So $SU(2) \times SU(2) \subset G_2$, $SU(3) \subset G_2$
(or $SO(4)$)

$$\begin{aligned} \bullet \hat{F}_4 & \text{ gives } F_4 \\ A_1 \times C_3 & \equiv SU(2) \times Sp(\quad) \\ A_2 \times A_2 & \equiv SU(3) \times SU(3) \\ A_3 \times A_1 & \equiv SU(4) \times SU(2) \\ B_4 & \equiv SO(9) \end{aligned}$$

$\hat{E}_8$ (also called E_9) gives

$$\begin{aligned}
 &E_8 \\
 &A_1 \times E_7 \\
 &A_2 \times E_6 \\
 &A_3 \times D_5 \equiv SU(4) \times SO(10) \equiv SO(6) \times SO(10) \\
 &A_4 \times A_4 = SU(5) \times SU(5) \quad (\text{because } A_3 = D_3) \\
 &A_1 \times A_2 \times A_5 \\
 &A_1 \times A_7 \\
 &D_8 \equiv SO(16) \\
 &A_8 \equiv SU(9)
 \end{aligned}$$

So $A_3 \times D_5 = SO(6) \times SO(10) \subset SO(16) = D_8$

So $A_3 \times D_5$ is not maximal (can embed it in another subgroup)

$[\alpha]$ [see p.4]

3.

[notes copied from
board p.5 → end here]

[notes copied from
board p.6 → begin here]

$$[2] E_8 \supset SU(5) \times SU(5)$$

$$248 \rightarrow \underbrace{(24, 1) + (1, 24)}_{\text{not enough}} + \underbrace{(5, \bar{5}) + (\bar{5}, 5)}_{\substack{\text{(funda, adj)} + (\\ \text{still not enough}}}$$

$$\text{or } + (5, \overline{10}) + (10, \bar{5}) \\ + (5, 10) + (\bar{5}, \overline{10}) \quad \text{OK?}$$

• $\hat{E}_6$ gives E_6

$$A_1 \times A_5 \quad \text{so } A_1 \times A_2 \times A_5 \text{ in } \hat{E}_8$$

$$A_2 \times A_2 \times A_2 \quad \left| \begin{array}{l} \text{not maximal because} \\ (A_1 \times A_5) \times A_2 \subset A_2 \times E_6 \subset E_8 \end{array} \right.$$

(and again the same)

For real rep of G of dim n , we can embed it in $SO(n)$

$$G \supset SO(n)$$

$$n \longleftarrow n$$

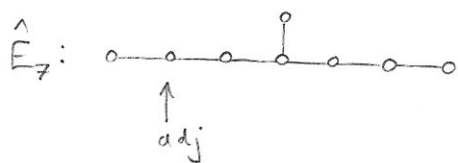
Example

$$SU(n) \subset SO(n^2 - 1)$$

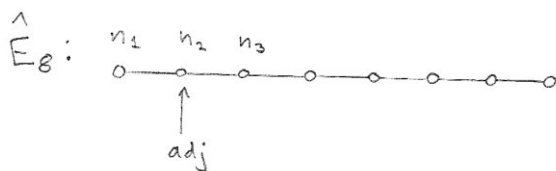
$$\text{adj} \longleftarrow (n^2 - 1) \text{ rep}$$

from the dim of the adj
to the rep of $SO(\dim \text{adj})$

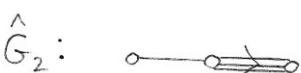
[Notes copied from
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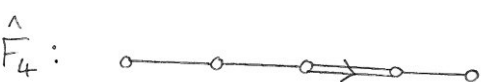
$\begin{bmatrix} & \\ & \end{bmatrix} [2, 0, 0, 0]$



$\begin{bmatrix} & \\ & \\ & \end{bmatrix} [0, 1, 0, 0]$

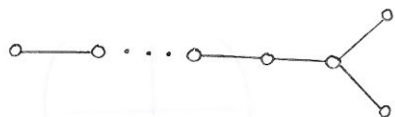


$\begin{bmatrix} \\ \\ \\ \end{bmatrix} [0, 0, 1, 0]$

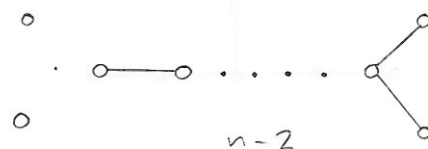
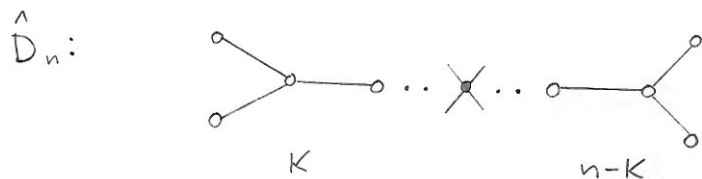


$A_n: [1, 0 \dots 0, 1]$

$B_n: [0, 1, 0 \dots 0]$



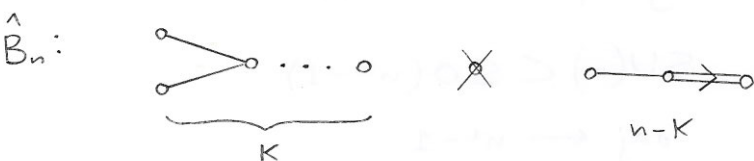
remove a node from the extended Dynkin diagram $\hat{G}$
remaining diagram is a subgroup of G



$SO(2n-4) \times SU(2) \times SU(2)$

$D_n \supset D_K \times D_{n-K}$

$SO(2n) \supset SO(2K) \times SO(2n-2K)$



$SO(2n+1) \supset SO(2K) \times SO(2n-2K+1)$

$\hat{G}_2 \supset G_2$

$SO(2n+1) \supset SO(2n)$

$\circ \dots \circ \supset A_1 \times A_1$
 $\supset A_2$

$SU(2) \times SU(2) \subset G_2$
 $SU(3) \subset G_2$

Branching Rules

[cf pp. 3-4]

[copied from
board after
lecture ended]

$$F_4, A_1 \times C_3, A_2 \times A_2, A_3 \times A_1, B_4$$

$$\hat{E}_8 = E_9$$

$$SO(6)$$

$$SU(4) \times SO(10)$$

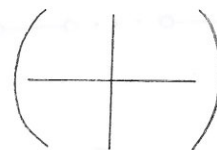
$$E_8, A_1 \times E_7, A_2 \times E_6, A_3 \times D_5, A_4 \times A_4$$

$$A_1 \times A_2 \times A_5, A_1 \times A_7, D_8, A_8 \rightarrow \text{not maximal}$$

not
maximal

$$SO(16) \subset E_8$$

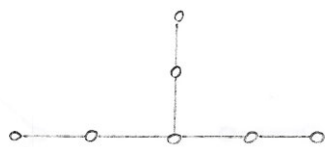
$$A_1 \times A_2 \times A_5 \subset A_2 \times E_6 \subset E_8$$

Dyn Kin

$$E_8 \supset SU(5) \times SU(5)$$

$$248 \rightarrow (24, 1) + (1, 24) + (5, \bar{10}) + (10, \bar{5}) + (5, 10) + (\bar{5}, \bar{10})$$

$$\hat{E}_6$$



$$\Lambda = \sum n_i \Lambda_i$$

$$\{n_i\}$$

$$E_6, A_1 \times A_5, A_2 \times A_2 \times A_2$$

$$\underline{SU(3)^3 \subset E_6}$$

For a real representation of a group G of dimension n

$$G \subset SO(n)$$

$$n \leftarrow n$$

Example

$$SU(n) \subset SO(n^2 - 1)$$

$$\text{adj} \leftarrow n^2 - 1$$